

AN EXAMPLE OF ONE - DIMENSIONAL PHASE TRANSITION

U.A. Rozikov

Institute of Mathematics, 29, F.Hodjaev str., 700143, Tashkent, Uzbekistan.

e-mail: rozikovu@yandex.ru

Abstract. In the paper a one-dimensional model with nearest - neighbor interactions $I_n, n \in \mathbf{Z}$ and spin values ± 1 is considered. We describe a condition on parameters I_n under which the phase transition occurs. In particular, we show that the phase transition occurs if $I_n \geq |n|, n \in \mathbf{Z}$.

Keywords: Configuration; One-dimension; Phase transition; Gibbs measure

1 Introduction

Interest in phase transition in one-dimensional systems has gained due to Little's work [8]. Existence or nonexistence depends markedly upon the model employed. In [11] Van Hove shows (see also [9, section 5.6.]) a one-dimensional system can not exhibit a phase transition if the (translation-invariant) forces are of finite range.

An interesting one-dimensional model considered by Baur and Nosanow [1], giving rise to a phase transition with only nearest- neighbor interactions which has some of the interaction constants equal to minus infinity. Note that in a particular case the model which we shall consider here is a model with only nearest-neighbor interactions without assumption of existence a interaction equal to minus infinity (cf. with [1]).

Other examples of one-dimensional phase transitions for models with long range interactions were considered in [3-7], [10].

In the paper we consider the Hamiltonian

$$H(\sigma) = \sum_{l=(x-1,x):x \in \mathbf{Z}} I_x \mathbf{1}_{\sigma(x-1) \neq \sigma(x)}, \quad (1)$$

where $\mathbf{Z} = \{\dots, -1, 0, 1, \dots\}$, $\sigma = \{\sigma(x) \in \{-1, 1\} : x \in \mathbf{Z}\} \in \mathcal{O} = \{-1, 1\}^{\mathbf{Z}}$, and $I_x \in R$ for any $x \in \mathbf{Z}$.

The goal of this paper is to describe a condition on parameters of the model (1) under which the phase transition occurs.

2 Phase transition

Let us consider the sequence $\mathbf{L}_n = [-n, n], n = 0, 1, \dots$ and denote $\mathbf{L}_n^c = \mathbf{Z} \setminus \mathbf{L}_n$. Consider boundary condition $\sigma_n^{(+)} = \sigma_{\mathbf{L}_n^c} = \{\sigma(x) = +1 : x \in \mathbf{L}_n^c\}$. The energy $H_n^+(\sigma)$ of the

configuration σ in the presence of boundary condition $\sigma_n^{(+)}$ is expressed by the formula

$$H_n^+(\sigma) = \sum_{l=(x-1,x):x \in \mathbb{L}_n} I_x \mathbf{1}_{\sigma(x-1) \neq \sigma(x)} + I_{-n} \mathbf{1}_{\sigma(-n) \neq 1} + I_{n+1} \mathbf{1}_{\sigma(n) \neq 1}. \quad (2)$$

The Gibbs measure on $\mathcal{O}_n = \{-1, 1\}^{\mathbb{L}_n}$ with boundary condition $\sigma_n^{(+)}$ is defined in the usual way

$$\mu_{n,\beta}^+(\sigma) = Z^{-1}(n, \beta, +) \exp(-\beta H_n^+(\sigma)), \quad (3)$$

where $\beta = T^{-1}$, $T > 0$ — temperature and $Z(n, \beta, +)$ is the normalizing factor (statistical sum).

Denote by σ_n^+ the configuration on \mathbf{Z} such that $\sigma_n^+(x) \equiv +1$ for any $x \in \mathbb{L}_n^c$.

Put

$$A(\sigma_n^+) = \{x \in \mathbf{Z} : \sigma_n^+(x) = -1\},$$

$$\partial(\sigma_n^+) = \{l = (m-1, m) \in \mathbf{Z} \times \mathbf{Z} : \sigma_n^+(m-1) \neq \sigma_n^+(m)\}.$$

Note that there is one-to-one correspondence between the set of all configurations σ_n^+ and the set of all subsets A of \mathbb{L}_n .

Let $A'(\sigma_n^+)$ be the set of all maximal connected subsets of $A(\sigma_n^+)$.

LEMMA 1. *Let $B \subset \mathbf{Z}$ be a fixed connected set and $p_\beta^+(B) = \mu_{n,\beta}^+\{\sigma_n^+ : B \in A'(\sigma_n^+)\}$. Then*

$$p_\beta^+(B) \leq \exp \left\{ -\beta \left[I_{n_B} + I_{N_B+1} \right] \right\},$$

where n_B (resp. N_B) is the left (resp. right) endpoint of B .

Proof. Denote $F_B = \{\sigma_n^+ : B \in A'(\sigma_n^+)\}$ — the set of all configurations σ_n^+ on Z with "+"-boundary condition (i.e. $\sigma_n^+(x) \equiv +1$ for any $x \in \mathbb{L}_n^c$) such that B is maximal connected set. Denote also $F_B^- = \{\sigma_n^+ : B \cap A'(\sigma_n^+) = \emptyset\}$. Define the map $\chi_B : F_B \rightarrow F_B^-$ as following: for $\sigma_n \in F_B$ we destroy the set B changing the values $\sigma_n(x)$ inside of B to $+1$. The constructed configuration is $\chi_B(\sigma_n) \in F_B^-$.

For a given B the map χ_B is one-to-one map.

For $\sigma_n \in F_B$ it is clear that

$$A'(\sigma_n) = A'(\chi_B(\sigma_n)) \cup B, \quad \partial(\sigma_n) = \partial(\chi_B(\sigma_n)) \cup \{(n_B-1, n_B), (N_B, N_B+1)\}.$$

Thus we have

$$H_n^+(\sigma_n) - H_n^+(\chi_B(\sigma_n)) = I_{n_B} + I_{N_B+1}. \quad (4)$$

By definition we have

$$p_\beta^+(B) = \frac{\sum_{\sigma_n \in F_B} \exp\{-\beta H_n^+(\sigma_n)\}}{\sum_{\sigma_n} \exp\{-\beta H_n^+(\sigma_n)\}} \leq \frac{\sum_{\sigma_n \in F_B} \exp\{-\beta H_n^+(\sigma_n)\}}{\sum_{\sigma_n \in F_B} \exp\{-\beta H_n^+(\chi_B(\sigma_n))\}}. \quad (5)$$

Using (4) from (5) we get

$$p_\beta^+(B) \leq \frac{\sum_{\sigma_n \in F_B} \exp \left\{ -\beta H_n^+(\chi_B(\sigma_n)) - \beta [I_{n_B} + I_{N_B+1}] \right\}}{\sum_{\sigma_n \in F_B} \exp \{ -\beta H_n^+(\chi_B(\sigma_n)) \}} = \exp \left\{ -\beta [I_{n_B} + I_{N_B+1}] \right\}.$$

The lemma is proved.

Assume that for any $r \in \{1, 2, \dots\}$ and $n \in \mathbf{Z}$ the Hamiltonian (1) satisfies the following condition

$$I_n + I_{n+r} \geq r. \quad (6)$$

LEMMA 2. Assume condition (6) is satisfied. Then for all sufficiently large β , there is a constant $C = C(\beta) > 0$, such that

$$\mu_\beta^+ \{ \sigma_n : |B| > C \ln |\mathbb{L}_n| \text{ for some } B \in A'(\sigma_n) \} \rightarrow 0, \text{ as } |\mathbb{L}_n| \rightarrow \infty,$$

where $|\cdot|$ denotes the number of elements.

Proof. Suppose $\beta > 1$, then by Lemma 1 and condition (6) we have

$$\mu_\beta^+ \{ \sigma_n : B \in A'(\sigma_n), t \in B, |B| = r \} = \sum_{B: t \in B, |B|=r} p_\beta^+(B) \leq r \exp \{ -\beta r \}.$$

Hence

$$\begin{aligned} \mu_\beta^+ \{ \sigma_n : B \in A'(\sigma_n), t \in B, |B| > C_1 \ln |\mathbb{L}_n| \} &\leq \sum_{r \geq C_1 \ln |\mathbb{L}_n|} r \exp \{ -\beta r \} \leq \\ &\sum_{r \geq C_1 \ln |\mathbb{L}_n|} \exp \{ (1 - \beta)r \} = \frac{|\mathbb{L}_n|^{C_1(1-\beta)}}{1 - e^{1-\beta}}, \end{aligned} \quad (7)$$

where C_1 will be defined latter. Thus we have

$$\mu_\beta^+ \{ \sigma_n : \exists B \in A'(\sigma_n), |B| > C_1 \ln |\mathbb{L}_n| \} \leq \frac{|\mathbb{L}_n|^{C_1(1-\beta)+1}}{1 - e^{1-\beta}}.$$

The last expression tends to zero if $|\mathbb{L}_n| \rightarrow \infty$ and $C_1 > \frac{1}{\beta-1}$. The lemma is proved.

LEMMA 3. Assume the condition (6) is satisfied. Then

$$\mu_\beta^+ \{ \sigma_n : \sigma_n(0) = -1 \} \rightarrow 0, \text{ as } \beta \rightarrow \infty. \quad (8)$$

Proof. If $\sigma_n(0) = -1$, then 0 is point for some $B \in A'(\sigma_n)$. Consequently,

$$\mu_\beta^+ \{ \sigma_n : 0 \in B, |B| < C_1 \ln |\mathbb{L}_n| \} \leq \sum_{r=1}^{C_1 \ln |\mathbb{L}_n|} (e^{1-\beta})^r \leq \frac{e^{1-\beta}}{1 - e^{1-\beta}}$$

and

$$\mu_{\beta}^{+}\{\sigma_n(0) = -1\} \leq \mu_{\beta}^{+}\{\sigma_n : 0 \in B, \quad B \in A'(\sigma_n)\} \leq \frac{e^{1-\beta}}{1 - e^{1-\beta}} + \frac{|\mathbf{L}_n|^{C_1(1-\beta)+1}}{1 - e^{1-\beta}}. \quad (9)$$

For $|\mathbf{L}_n| \rightarrow \infty$ and $\beta \rightarrow \infty$ from (9) we get (8). The lemma is proved.

THEOREM 4. *Assume the condition (6) is satisfied. For all sufficiently large β there are at least two Gibbs measures for the model (1) .*

Proof. Using a similar argument one can prove

$$\mu_{\beta}^{-}\{\sigma_n : \sigma_n(0) = 1\} \rightarrow 0, \quad \text{as } \beta \rightarrow \infty.$$

Consequently, for sufficiently large β we have

$$\mu_{\beta}^{+}\{\sigma_n : \sigma_n(0) = -1\} \neq \mu_{\beta}^{-}\{\sigma_n : \sigma_n(0) = -1\}.$$

This completes the proof.

Denote

$$\mathcal{H} = \{H : H \text{ (see (1)) satisfies the condition (6)}\}$$

The following example shows that the set \mathcal{H} is not empty.

Example. Consider Hamiltonian (1) with $I_m \geq |m|$, $m \in \mathbf{Z}$. Then

$$I_m + I_{m+k} \geq |m| + |m+k| \geq k$$

for all $m \in \mathbf{Z}$ and $k \geq 1$. Thus the condition (6) is satisfied.

Acknowledgments. The work supported by NATO Reintegration Grant : FEL. RIG. 980771. The final part of this work was done in the Physics Department of “La Sapienza” University in Rome. I thank M.Cassandro and G.Gallavotti for an invitation to the “La Sapienza” University and useful discussions. I thank the referee for useful suggestions.

References

1. Baur M., Nosanow L. *J. Chemical Phys.* **37**: 153-160 (1962).
2. Cassandro M., Ferrari P., Merola I., Presutti E. *arXiv: math-ph/ 021 1062*, **2** 28 Nov 2002.
3. Dyson F. *Commun. Math. Phys.* **12**: 91-107 (1969)
4. Dyson F. *Commun. Math. Phys.* **21**: 269-283 (1971)
5. Flohlich J., Spencer T. *Commun. Math. Phys.* **84**: 87-101 (1982)
6. Johansson K. *Commun. Math. Phys.* **141**: 41-61 (1991)
7. Johansson K. *Commun. Math. Phys.* **169** : 521-561 (1995)
8. Little W., *Phys. Rev.* **134**: A1416-A1424 (1964)
9. Ruelle D., *Statistical mechanics (rigorous results)*, Benjamin, New York, 1969
10. Strecker J. *J. Math. Phys.* **10**: 1541-1554 (1969)
11. Van Hove L. *Physica* **16**: 137-143 (1950).